Bayesian forward modeling of high-resolution radio interferometric gravitational lens observations



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Motivation

Science with strong lenses

Detection of low-mass dark matter haloes, constraints on mass function

e.g. Ritondale et al. (2019), Despali et al. (2018), Birrer et al. (2017), Hezaveh et al. (2016), Vegetti et al. (2014), Vegetti et al. (2012), Vegetti et al. (2010)

• Source science taking advantage of magnification

e.g. Spingola et al. (2019), Rizzo et al. (2018), Johnson et al. (2017), Leethochawalit et al. (2016), . Swinbank et al. (2015)

Cosmology using time delays between images

e.g. Rusu et al. (2019), Suyu et al. (2018), Wong et al. (2017), Treu and Marshall (2016), Chen et al. (2016), Courbin et al. (2011), Fassnacht et al. (2002)

Polarimetry and rotation measure synthesis

e.g. Mao et al. (2017)

Lots of science with lenses - review in Koopmans et al. (2009).

Science with strong lenses

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Mass function: CDM vs. WDM



Lovell et al. (2014)

Mass function: CDM vs. WDM



Despali and Vegetti (2017), Lovell et al. (2014)

LoW-Mass Perturbers (LwMPs) with lensing







LoW-Mass Perturbers (LwMPs) with lensing







Global VLBI



- GVLBI array (Global Very Long Baseline Interferometry)
- Sub-milliarcsec resolution at shorter wavelengths
- Huge numbers of visibilities (up to $\sim 10^{10}).$ A challenge for modeling.

Modeling

Forward modeling



ALMA (ESO/NRAO/NAOJ), L. Calada (ESO), Y. Hezaveh et al.

$$P(\mathbf{s} \mid \mathbf{d}, \eta, \lambda) = \frac{P(\mathbf{d} \mid \mathbf{s}, \eta) P(\mathbf{s} \mid \lambda)}{P(\mathbf{d} \mid \eta, \lambda)}$$

- d is the data (complex radio visibilities)
- s is the source plane surface brightness
- η is a vector containing lens parameters (mass, slope, ellipticity, LwMPs...)
- λ is a *hyper-parameter* setting the strength of the regularization (fixed for MAP)

Multi-level (hierarchical) Bayes

We use a three-step inference process:

1. Optimize for linear parameters (source inversion)

$$\mathsf{Max}\left[P(\mathbf{s} \mid \mathbf{d}, \eta, \lambda) = \frac{P(\mathbf{d} \mid \mathbf{s}, \eta) P(\mathbf{s} \mid \lambda)}{P(\mathbf{d} \mid \eta, \lambda)}\right]$$

2. Optimize for non-linear parameters (lens model and regularization)

$$\mathsf{Max} \left| P(\eta, \lambda \mid \mathbf{d}) = \frac{P(\mathbf{d} \mid \eta, \lambda) P(\eta, \lambda)}{P(\mathbf{d})} \right|$$

3. Evidence computation and model comparison

$$\mathsf{Max} \left| P(\mathbf{d}) = \int P(\mathbf{d}|\eta, \lambda) P(\eta, \lambda) \, \mathrm{d}\lambda \, \mathrm{d}\eta \right|$$

See Vegetti and Koopmans (2009) for the specific implementation, MacKay (1991) for a theoretical overview.

Forward modeling



ALMA (ESO/NRAO/NAOJ), L. Calada (ESO), Y. Hezaveh et al.

Maximizing the log-posterior with respect to *s* implies

$$\frac{\partial}{\partial s} \left[(D L s - d)^T C_d^{-1} (D L s - d) + s^T R s \right] = 0$$
giving
$$(L^T D^T C_d^{-1} D L + R) s = L^T D^T C_d^{-1} d$$

$$(L^{T}D^{T}C_{d}^{-1}DL + R)s = L^{T}D^{T}C_{d}^{-1}d$$

Vectors:

- s Reconstructed source solution
- *d* Data (radio visibilities)

Matrices:

- L Lensing operator
- C_d^{-1} Visibility-space noise covariance
- R Regularization matrix
- D Nonuniform discrete Fourier transform

Replace discrete FT with non-uniform FFT:

 $D \to G F W$

- G Gridding operator
- *F* FFT
- W Apodization operator (inverse FT of the gridding kernel)

$$(L^T W^T F^T G^T C_d^{-1} G F W L + R) s = L^T W^T F^T G^T C_d^{-1} d$$

Vectors:

- s Reconstructed source solution
- *d* Data (radio visibilities)

Explicit matrices:

- L Lensing operator
- C_d^{-1} Visibility-space noise covariance
- *R* Regularization matrix
- W Apodization operator (inverse FT of the gridding kernel)

$$(L^T W^T F^T G^T C_d^{-1} G F W L + R) s = L^T W^T F^T G^T C_d^{-1} d$$

Implicitly defined linear operators ("Matrix-Free"):

- *G* Visibility gridding operator (implemented on GPU for speed)
- *F* FFT

The Fourier Transform is a dense matrix. Large numbers of visibilites (> 10^8) demand that we use a gridding operation + FFT instead of direct DFT.

For space, denote:

$$A \equiv (L^T W^T F^T G^T C_d^{-1} G F W L + R)$$

and

$$b \equiv L^T W^T F^T G^T C_d^{-1} d$$

- A is not explicitly defined. It exists only as a function for matrix-vector multiplication (linearly mapping vectors to vectors)
- We must solve A s = b
- We must also compute log det A

Since A is not explicitly defined, we must use an iterative solver. Conjugate gradient seems to work best for this problem. However, we *must* precondition to get reasonably fast convergence, due to extremely high condition number (can be 10^{10}).

Instead of A s = b, we solve

$$P^{-1}As = P^{-1}b$$

We must design the preconditioner P such that it is a sparse matrix which approximates A reasonably well.

Preconditioned conjugate gradient solver

Our preconditioner is based on the image-plane noise covariance:

$$C_x^{-1} = W^T F^T G^T C_d^{-1} G F W$$

One can show that each row of this matrix simply contains the naturally-weighted dirty beam. We approximate

$$\tilde{C}_x^{-1} = \operatorname{diag}(C_x^{-1})$$

which is a diagonal matrix containing the brightest pixel in the dirty beam. Using only the diagonal works well and guarantees positive-definiteness of P.

Our full preconditioner is then

$$P = L^T \tilde{C}_x^{-1} L + R$$

which is sparse and can be decomposed for fast application of P^{-1} . 20

For space, denote:

$$A \equiv (L^T W^T F^T G^T C_d^{-1} G F W L + R)$$

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In optimizing the lensing operator L, we compute the Bayesian evidence. Correct normalization of this evidence requires us to compute logdet(A).

We find that the determinant of our sparse preconditioner also provides an adequate approximation for this normalization factor.

$$\log\det(A) \approx \log\det P = \log\det(L^T \tilde{C}_x^{-1} L + R)$$

The determinant is computed for free when performing the Cholesky decomposition.

Results

Kinematics of SPT0418-47 (Francesca Rizzo)



Dusty starburst galaxy at z = 4.2, effective ~ 60 pc resolution. First study of $v_{\rm rot}/\sigma$ at this redshift.

Kinematics of SPT0418-47 (Francesca Rizzo)



Rizzo et al., "A dynamically cold disk galaxy in the early Universe", *Nature* 2020.



Y. Tamura (The University of Tokyo)/ALMA (ESO/NAOJ/NRAO)



SDP.81 Band 6+7 observation. Attempting to reproduce substructure detection of Hezaveh et al. (2016).

Computed from 8.8×10^8 visibilities. ~ 7 mas resolution.

SDP.81



Posterior sampling in \sim 12 hours on 100 cores + 16 GPUs.

MG J0751+2716



GVLBI data at \sim 5 mas resolution, component-based modeling on CLEANed images (shown here) done by Spingola et al. (2018).



 4×10^8 visibilities, 5 mas resolution. Cored elliptical power law lens model does not quite explain the data. Modeling ongoing...

Outlook

- Working forward modeling code for radio interferometric data
- GPU acceleration for visibility gridding/de-gridding operation
- Data distributed over global MPI communicator for efficient memory use
- Currently applying code to several lens systems for DM searches and source science (Devon, Simona, Francesca)

- Self-calibration
- Polarimetry and rotation measure synthesis
- Plasma lens effects: Time delay, plasma lensing

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